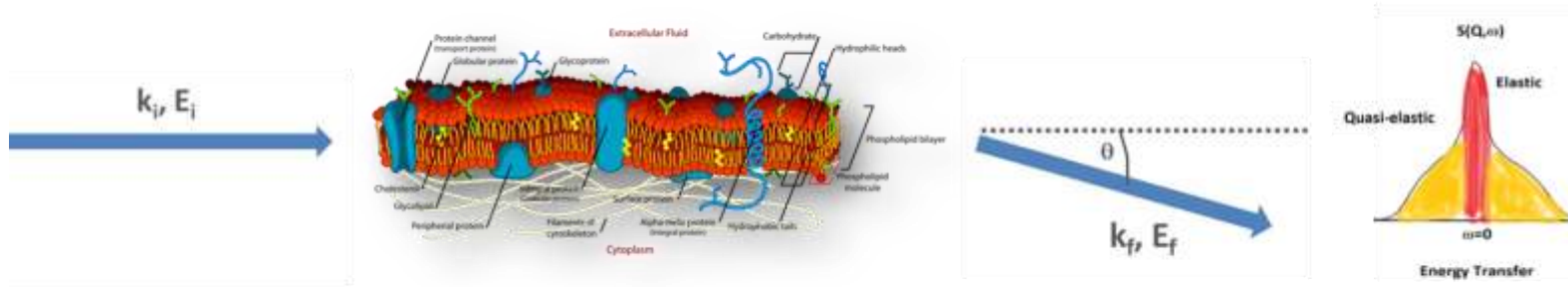


Dynamics of biomolecules from lipids to proteins



Victoria Garcia Sakai

ISIS Neutron and Muon Facility

Rutherford Appleton Lab, Didcot, UK



Science & Technology Facilities Council

ISIS

Outline

- Why** do we need to study nanoscale motions in biomolecules
- What is **quasi-elastic neutron scattering**
- What information do **neutrons** give us
- What **instruments** are available
- Examples**

Take home messages

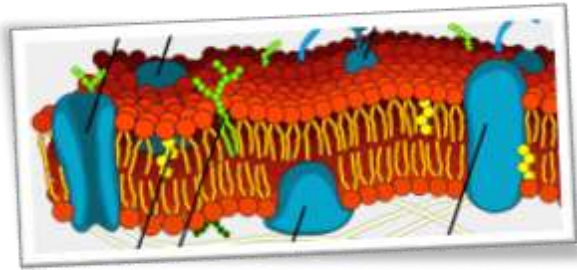
1. **Dynamics** provide insights into relevant biomolecules
2. **Instruments** available to probe different motions
3. **Complementary** to structural characterisation
4. Neutrons provide **unique** information

Importance of Dynamics for Biomolecules

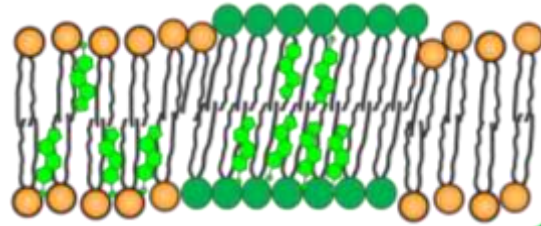
Function and Transport

- Enzymatic activity (internal motions, solvent dynamics)
- Ligand-binding and recognition (protein flexibility)
- Transport and permeability of lipids (dynamics of lipids and drugs)
- Lipid-protein, protein-protein interactions (membrane fluidity)
- Evolutionary understanding (protein folding and unfolding)
- Characteristics of extremophiles (protein flexibility, concentration dependence, and P and T dependence)
- Stability and endurance (rigidity)
- Drug delivery vectors (lipid dynamics)
- Diseases - cancer, Alzheimer's (lipid, protein and solvent dynamics)
- Cryopreservation (protein and solvent dynamics)
- ...

Protein-Protein and Protein-Membrane Interactions

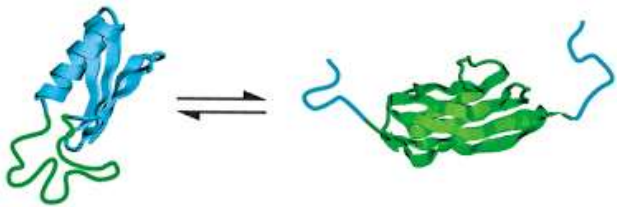


Lipids as model membranes

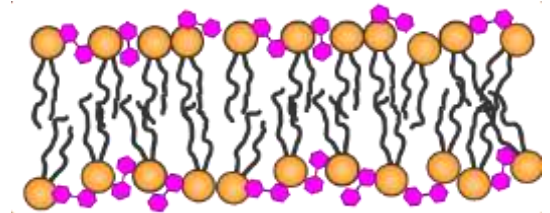


Effect of small molecules (drugs, carbohydrates ...)

Protein folding/unfolding



Fluidity
Permeability
Flexibility

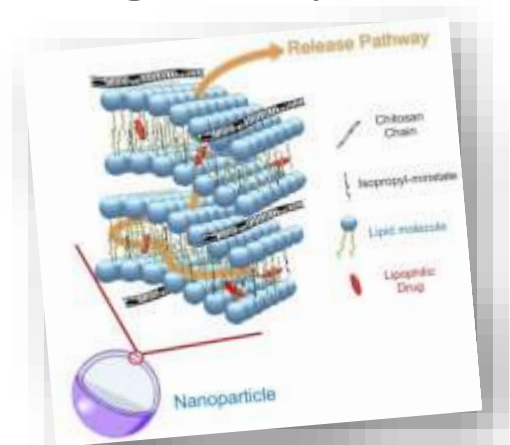
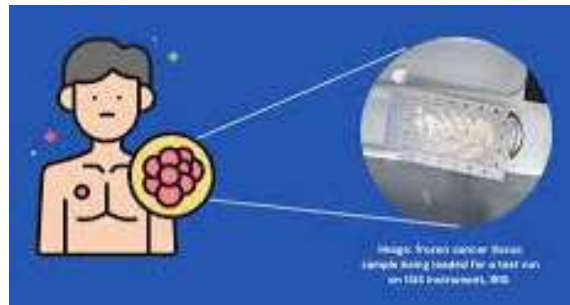


Design of drug delivery vectors

Understanding extremophiles



Cancer drug design



Neutrons for Studying Motions

The Nobel Prize in Physics 1994



Neutrons behave as particles and as waves

The Royal Swedish Academy of Sciences has awarded the 1994 Nobel Prize in Physics for pioneering contributions to the development of neutron scattering techniques for studies of condensed matter.

Clifford G. Shull, JPL, Cambridge, Massachusetts, USA, received one half of the 1994 Nobel Prize in Physics for development of the neutron diffraction technique.



S Shull made use of atoms scattering X-rays of neutrons which change direction without losing energy when they collide with atoms. Because of the wave nature of neutrons, a diffraction pattern can be recorded which indicates where in the sample the atoms are located. Even the placing of light electrons such as hydrogen in molecular hydroides, or hydrogen, carbon and nitrogen in organic substances can be determined. The pattern also shows how atoms vibrate as they vibrate in magnetic materials, since neutrons are affected by magnetic forces. Shull also made use of this phenomenon in his neutron diffraction technique.



Neutrons see neutrons, unlike X-rays
In contrast to X-rays, neutrons interact with neutrons. This means that they can see the neutrons in a sample, which is why they are used to study the structure of materials. Neutrons are also used to study the structure of materials, which is why they are used to study the structure of materials.

Neutrons reveal inner stresses
Neutrons can be used to study the inner stresses in a material. This is because neutrons interact with the nuclei of atoms, which are affected by the stresses in the material. This means that neutrons can be used to study the inner stresses in a material, which is why they are used to study the inner stresses in a material.

Neutrons show what atoms vibrate
Neutrons can be used to study what atoms vibrate in a material. This is because neutrons interact with the nuclei of atoms, which are affected by the vibrations in the material. This means that neutrons can be used to study what atoms vibrate in a material, which is why they are used to study what atoms vibrate in a material.

How it worked
Shull and Bragg made their pioneering contributions to the first neutron scattering facilities in the USA and Canada in the late 1940s and 1950s. It was then that the neutron scattering technique was developed further under the guidance of the neutron scattering community.

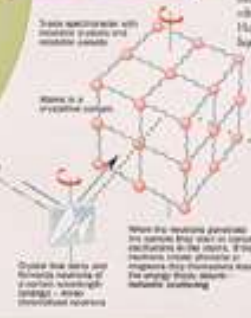
Neutrons reveal structure and dynamics

Neutrons show where atoms are



Neutrons bounce against atomic nuclei. They also react to the magnetism of the atoms.

Neutrons show what atoms do



Bertil H. Bragg, JPL, Cambridge, Massachusetts, USA, received one half of the 1994 Nobel Prize in Physics for development of neutron scattering techniques for studies of condensed matter.



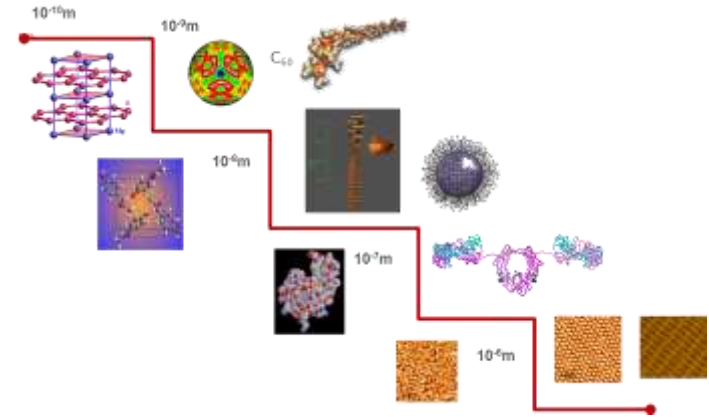
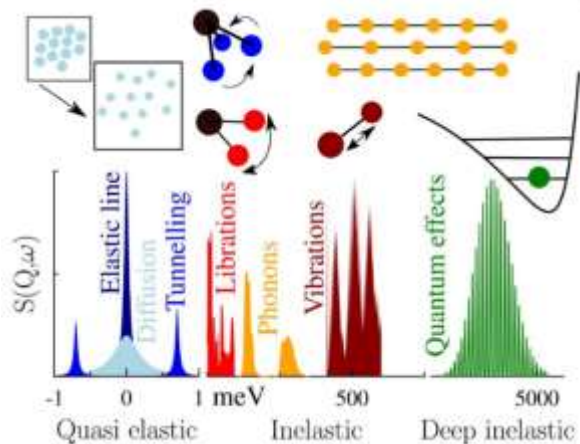
B Bragg made use of neutron scattering, i.e. of neutrons, which change their direction and energy when they collide with atoms. They then set up a series of atomic oscillations in crystals and social interactions in liquids and solids. Neutrons can also interact with light waves in magnets. With his 3-axis spectrometer Bragg measured energies of phonons (atomic vibrations) and magnons (magnetic waves). He also studied how atomic structures in liquids change with time.



Neutrons tell us where atoms are and what atoms do

Neutron Properties

1. Neutron wavelengths are of the same order as atomic spacings - **structural information** from Å to mm with structural precision to 10^{-15} m.



2. Neutron energies are of the same order as molecular modes – **dynamical information** from fs to ms and with spatial info (vibrations, diffusive motions).

Complementary to many other techniques:

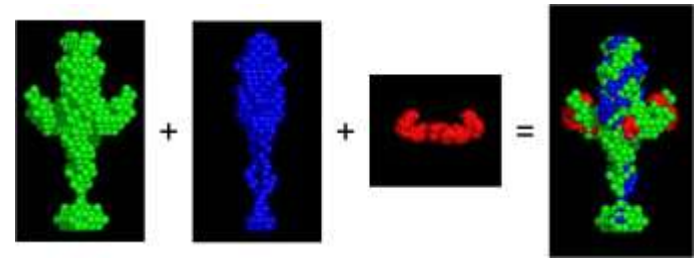
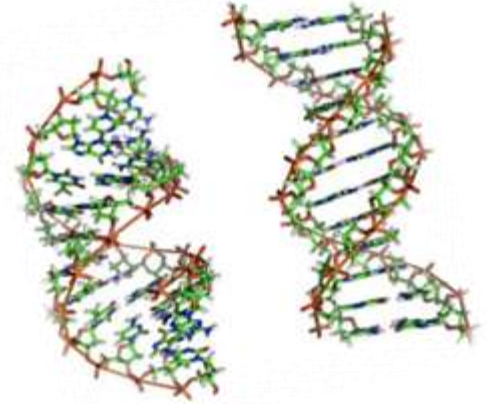
- X-rays
- Raman, IR
- Dynamic Light Scattering
- Photon Correlation Spectroscopy
- NMR
- MD simulations

Neutron Properties

3. Neutrons are **neutral** – highly penetrating, non-destructive.

4. Neutrons interact with nuclei – **sensitive to** light atoms, particularly ^1H , allows exploitation of isotopic variation (**H/D**).

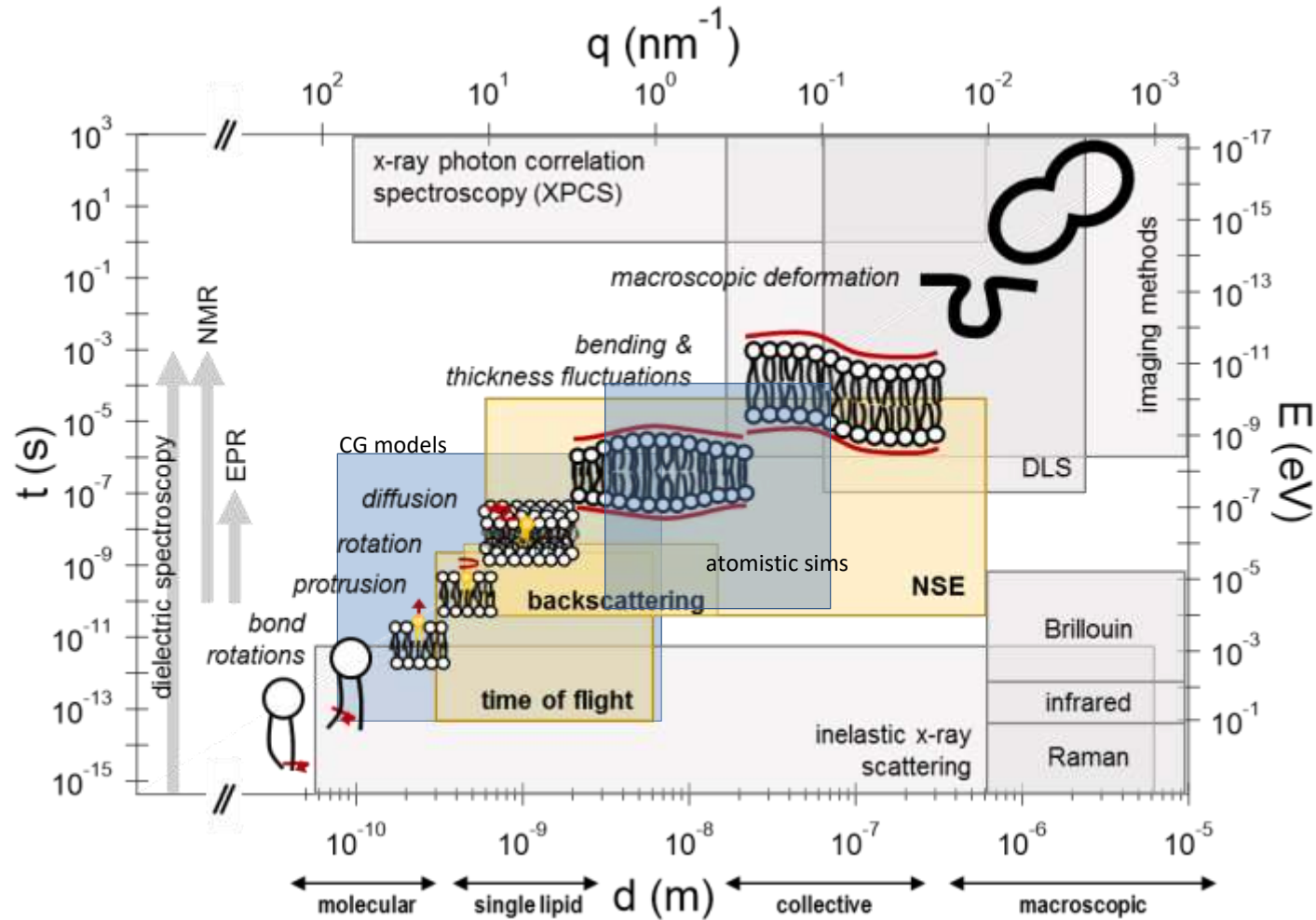
5. Neutrons have a **magnetic** moment—provides extra contrast and signal separation techniques.



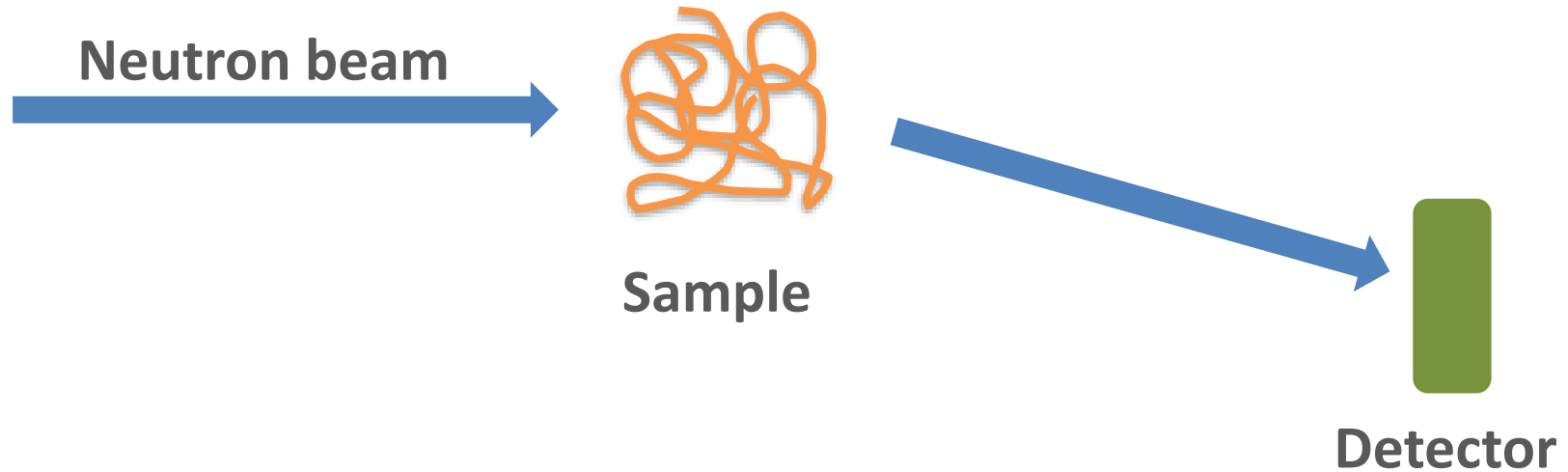
Unfortunately there are some **impracticalities**:

- neutrons are scarce and not available in your lab
- sample mass in some cases is large
- deuteration (isotopic substitution) is not always straight forward

Dynamical Modes for Lipid Membranes

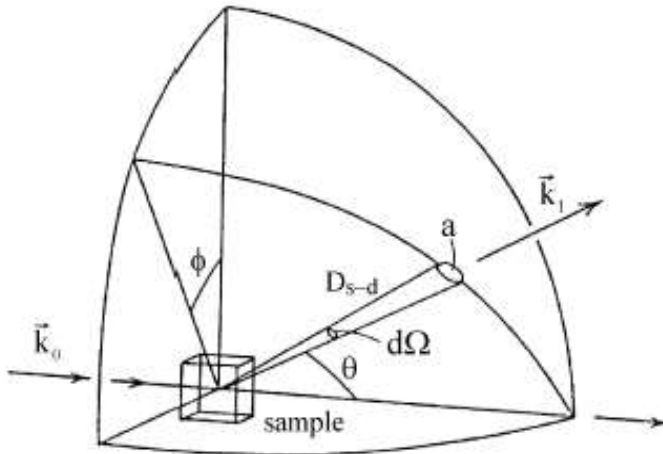


Neutron Scattering



- Source of neutron **radiation** on a sample
- An interesting sample
- Particles are scattered by the sample – this depends on the **interaction potential** between the two
- A neutron detector measures the **distribution** of radiation scattered
- A description of the measured signal in terms of the physics in the sample

The Scattering Function, $S(\mathbf{Q}, \omega)$



$$\frac{d^2\sigma}{d\Omega d\omega} \propto \frac{k_f}{k_i} \frac{\sigma}{4\pi} S(\mathbf{Q}, \omega)$$

The **scattering function, $S(\mathbf{Q}, \omega)$** contains all the physics of the system (in space and time) and depends only on the system.

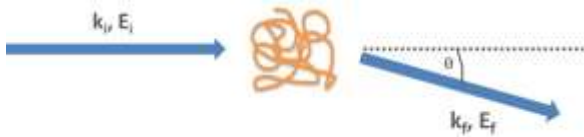
If your detector can analyse the energy of the neutrons, then the **double differential cross-section** can be defined as

$$\frac{d^2\sigma}{d\Omega dE} = \frac{\text{number of particles scattered per } s \text{ into a solid angle } d\Omega \text{ with final energies between } E_f \text{ and } (E_f + dE_f)}{I_0 d\Omega dE_f}$$

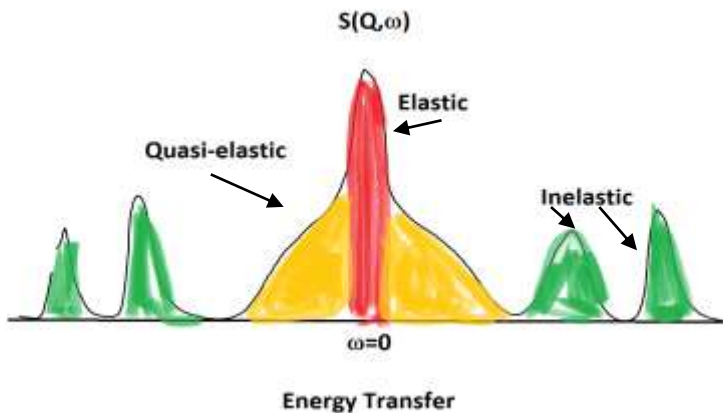
The Scattering Function, $S(Q, \omega)$

Information on dynamics

$$S(Q, \omega) = S(Q, \omega)_{\text{elastic}} + S(Q, \omega)_{\text{inelastic}} + S(Q, \omega)_{\text{quasi-elastic}}$$



Elastic scattering – no energy exchange $\hbar\omega=0$. Ideally it is a δ function, in reality it is the resolution.



Quasi-elastic scattering (QENS)– a small energy exchange $\hbar\omega \neq 0 \approx \text{neV}$ or few meV . Processes with a distribution of energies.

Inelastic scattering (INS) – energy exchange $\hbar\omega \neq 0$. Processes of discrete energy steps, quantised (vibrations or excitations).

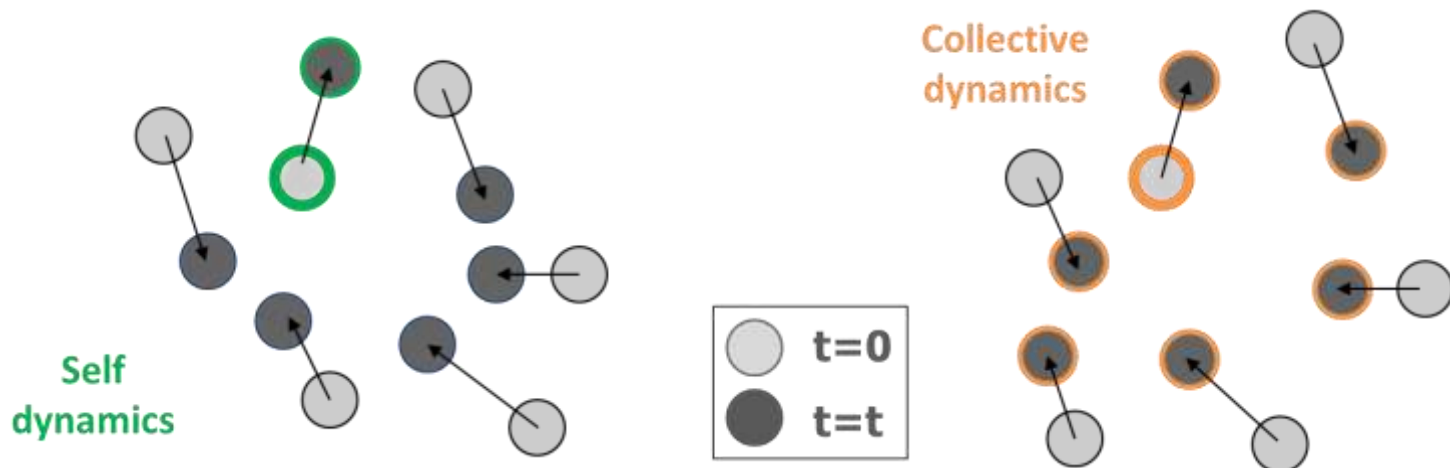
Scattering Function- Correlations

In the experiment we measure the total $S(\mathbf{Q}, \omega)$ and that each term, coherent and incoherent is weighted by its respective cross-section σ

$$S(\mathbf{Q}, \omega) = S_{\text{inc}}(\mathbf{Q}, \omega) + S_{\text{coh}}(\mathbf{Q}, \omega)$$

$$S_{\text{inc}}(\mathbf{Q}, \omega) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \sum_i \langle \exp(-i\mathbf{Q} \cdot \mathbf{R}_i(0)) \exp(-i\mathbf{Q} \cdot \mathbf{R}_i(t)) \rangle \exp(-i\omega t) dt$$

$$S_{\text{coh}}(\mathbf{Q}, \omega) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \sum_{i,j} \langle \exp(-i\mathbf{Q} \cdot \mathbf{R}_i(0)) \exp(-i\mathbf{Q} \cdot \mathbf{R}_j(t)) \rangle \exp(-i\omega t) dt$$



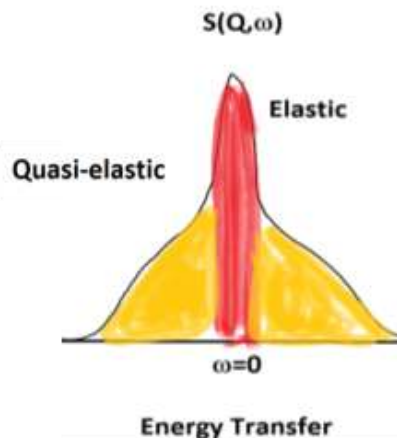
Dynamic Structure Factor, $S(Q, \omega)$

Incoherent scattering ...

- ... describes dynamics of **individual** particles
- ... molecular and single amphiphilic molecule
 - ... rotations, *trans-gauche* conformational transitions, diffusion
- ... main association of QENS technique (ps, ns)
 - ... relies on large incoh x-section of H*

Coherent scattering...

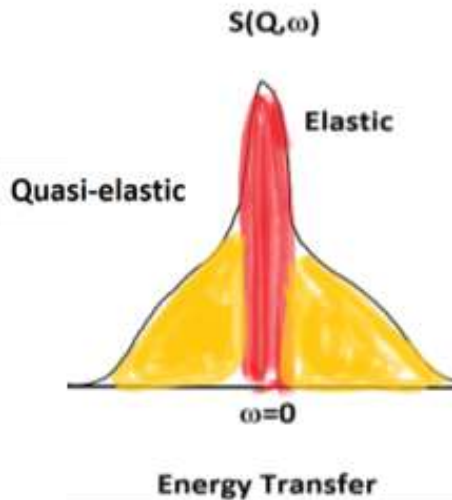
- ... describes **collective** dynamics of nuclei
- ... describes correlations between nuclei
 - ... bending and thickness fluctuations
 - ... strongly related to structure*
- ... main association with NSE technique (ns)



* Measurable is a function of Q and if coherent is considerable % of total signal need to pay attention to elastic structure factor $S(Q)$

Self-dynamics - Incoherent

Self-dynamics Incoherent QENS



We measure **the self correlation function**, ie. how particles move as a function of time. This corresponds to the incoherent signal which we can write as:

$$S_{\text{inc}}(\mathbf{Q}, \omega) = S_{\text{vib}}(\mathbf{Q}, \omega) \otimes S_{\text{rot}}(\mathbf{Q}, \omega) \otimes S_{\text{trans}}(\mathbf{Q}, \omega)$$

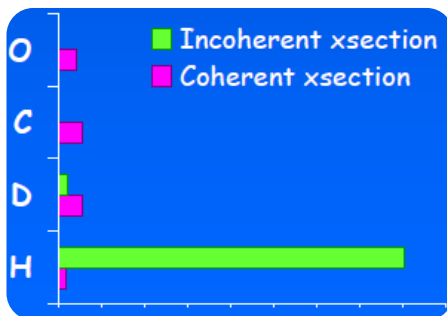
$$I_{\text{self}}(\mathbf{Q}, t) = I_{\text{vib}}(\mathbf{Q}, t) \times I_{\text{rot}}(\mathbf{Q}, t) \times I_{\text{trans}}(\mathbf{Q}, t)$$

It is a **convolution of components** which for simplicity are assumed **independent** motions. Note that in the time domain we multiply the terms (easier!)

Can't forget the instrumental resolution:

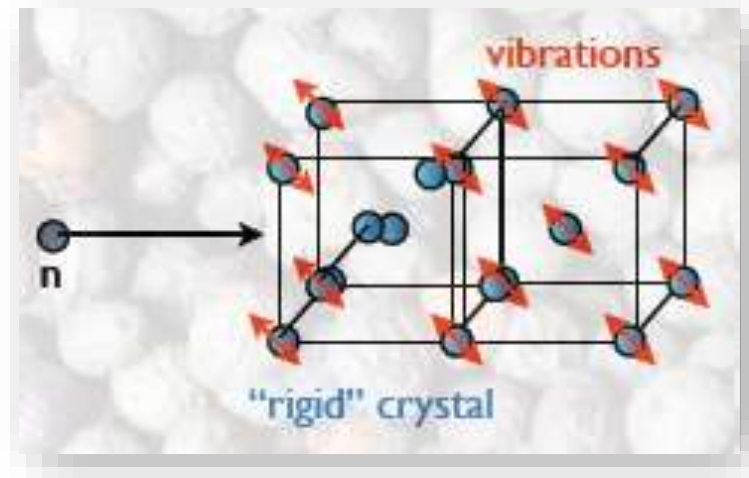
$$S'(\mathbf{Q}, \omega) = S(\mathbf{Q}, \omega) \otimes R(\mathbf{Q}, \omega)$$

QENS relies on the large incoherent xsection of hydrogen.



Vibrations: Debye-Waller factor

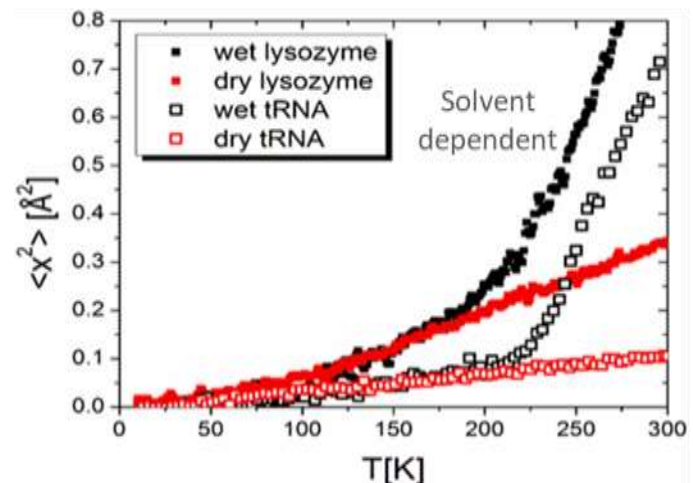
Nuclei (even in a rigid crystal) are **not stationary**, causing a decrease in the intensity of a diffracted beam (remember Bragg's law) because waves are not so well in phase. In addition, this smearing is a function of temperature. Assuming that **vibrations are harmonic and isotropic**.



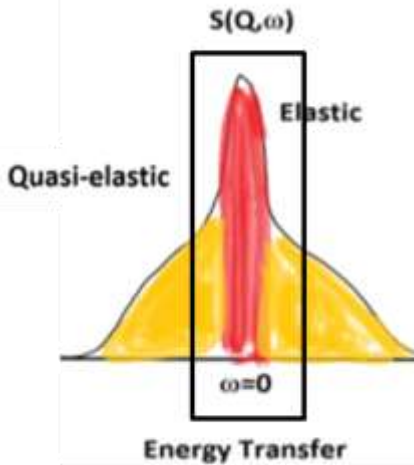
$$DWF = \langle \exp(i\mathbf{Q} \cdot \mathbf{u}) \rangle = \exp(-\langle (\mathbf{Q} \cdot \mathbf{u})^2 \rangle) = \frac{1}{3} \exp(Q^2 \langle u^2(T) \rangle)$$

Mean-square displacement of atoms, $\langle u^2 \rangle$, is measured as a function of temperature.

We can obtain an **effective force constant** –measure of flexibility



Fixed Energy Window Scans

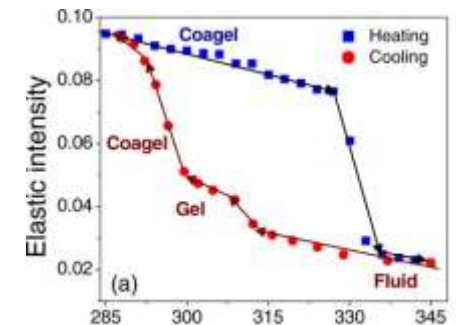
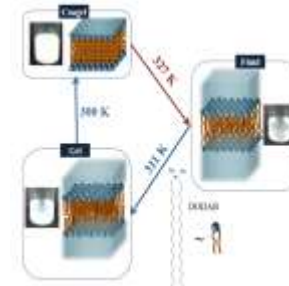
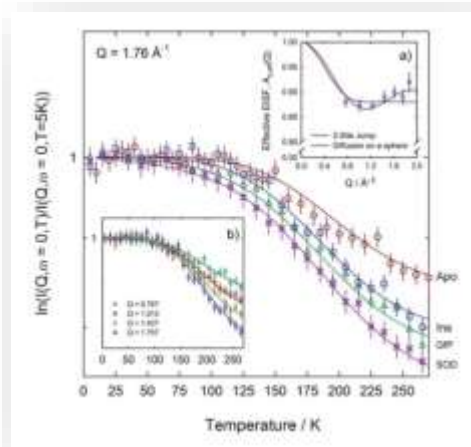


Case of $S(Q, \omega \approx 0)$: Elastic Fixed Window Scans

Measure the elastic intensity as a function of temperature (resembles a DSC scan).

Used to calculate $\langle u^2 \rangle$ using Gaussian approximation.

However, also good for locating **transitions**, at what temperatures do the **dynamics enter the time window**, and comparative studies (parametric).

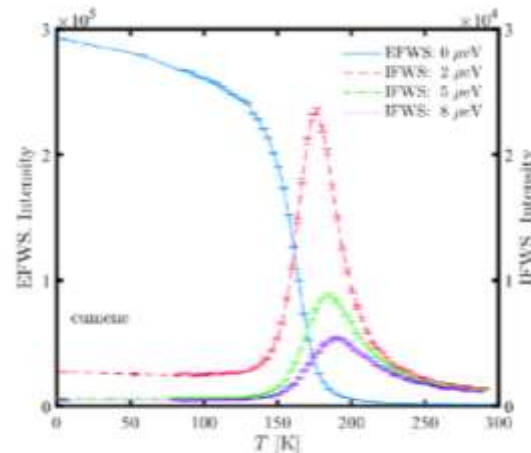
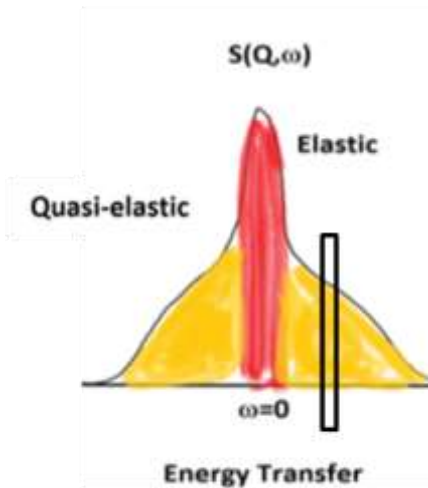


$$\frac{S^{\text{inc}}(Q, \omega=0, T)}{S^{\text{inc}}(Q, \omega=0, T \approx 0)} \propto \exp\left(-\frac{1}{3}Q^2 \langle r^2 \rangle\right) \text{ for } Q^2 \langle r^2 \rangle \leq 1$$

Fixed Energy Window Scans

Case of $S(Q, \omega \approx \text{offset})$: Inelastic Fixed Window Scans

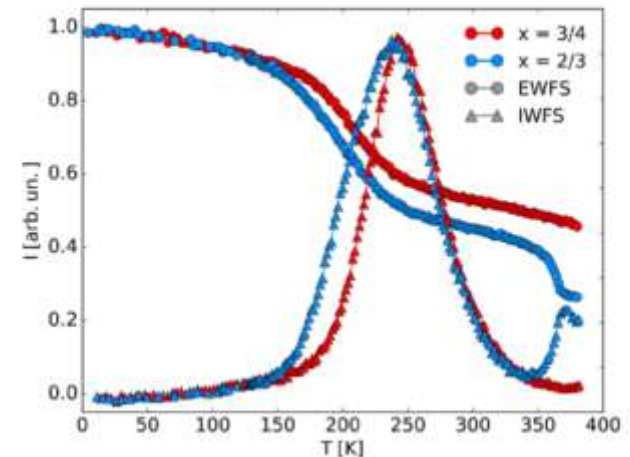
If transitions are not so clear or if difficult to gauge timescale, measure 'inelastic' intensity as a function of temperature. Slightly different measurement in reactor and spallation instruments.



Hansen et al, PRB, 2017

Different offsets correspond to different timescales

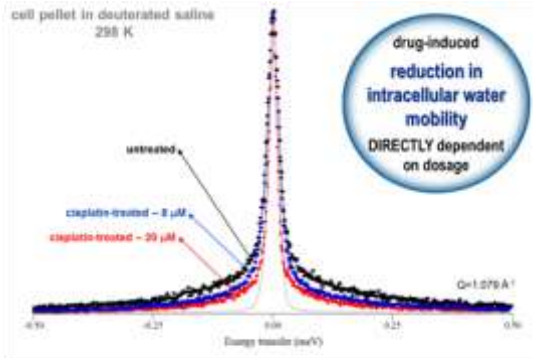
Peak shape and changes with Q provide additional information



Burankova et al, JPhysChemC, 2017

Diffusion

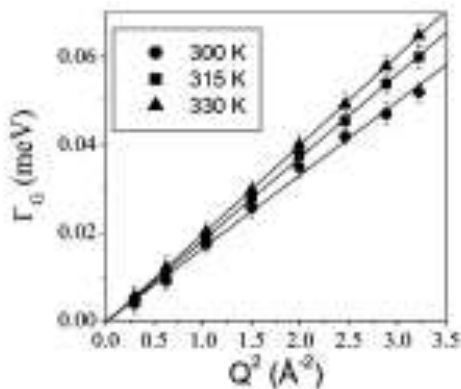
$$S_{\text{inc}}(\mathbf{Q}, \omega) = S_{\text{vib}}(\mathbf{Q}, \omega) \otimes S_{\text{rot}}(\mathbf{Q}, \omega) \otimes S_{\text{trans}}(\mathbf{Q}, \omega)$$



Batista de Carvalho et al, Phys Chem Chem Phys 19, 2017

QENS spectra describe by Lorentzians assigned to given diffusion model.

E.g. Fick's Law with Γ , the **half width at half maximum** is $= DQ^2$. Diffusion coefficient follows Arrhenius Law: $D \approx \exp(-Ea/kT)$.



- Diffusion coefficients of different motions
- Activation energies
- Lengths: confinement/jump
- Residence times (jumps)
- Geometry EISF ...

More models including rotations

$$S_{\text{inc}}(\mathbf{Q}, \omega) = S_{\text{vib}}(\mathbf{Q}, \omega) \otimes S_{\text{rot}}(\mathbf{Q}, \omega) \otimes S_{\text{trans}}(\mathbf{Q}, \omega)$$

**Elastic stationary
part, EISF**

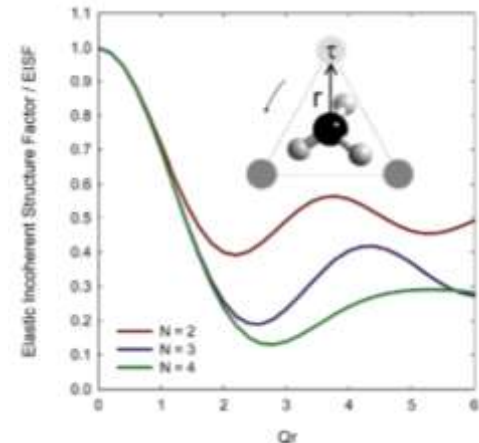
**Quasi-elastic
decaying part**

$$S_{\text{inc}}(\mathbf{Q}, \omega) = \exp(-Q^2 \langle u^2 \rangle) [A_0(\mathbf{Q}) \delta(\omega) + (1 - A_0(\mathbf{Q})) L(\mathbf{Q}, \omega)]$$

Localised motions including rotations give to a Q-independent line-width and a stationary part called the **Elastic Incoherent Structure Factor (EISF)** -fraction of elastic contribution. Gives information on the **geometry** of the motion/rotation.

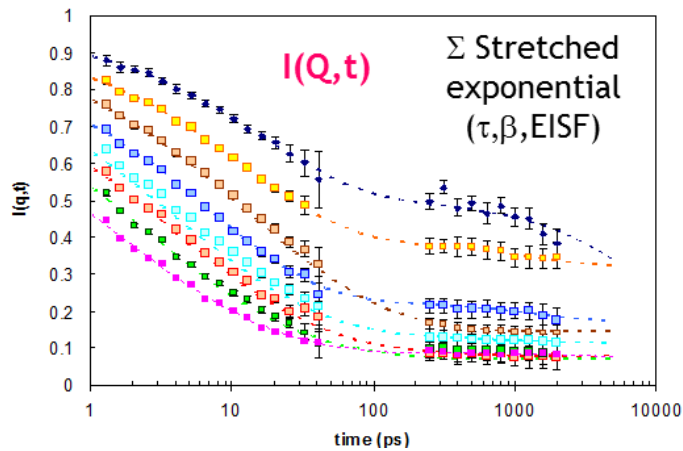
$$EISF = \frac{S_{\text{inc}}^{\text{el}}(\mathbf{Q})}{S_{\text{inc}}^{\text{el}}(\mathbf{Q}) + S_{\text{inc}}^{\text{qel}}(\mathbf{Q})}$$

- Jumps between 2, 3, ... n sites
- Rotational diffusion on a circle
- Diffusion on a sphere
- Diffusion inside a sphere, cylinder...



or characteristic relaxation

$I_{\text{self}}(\mathbf{Q}, t)$ = stretched exponential



Data is Fourier Transformed if necessary

Fit with one or more exponentials
(typically stretched)

Each one assigned to a specific dynamical mode

$$\frac{I(\mathbf{Q}, t)}{I(\mathbf{Q}, 0)} = A \exp \left[- \left(\frac{t}{\tau_{\text{KWW}}} \right)^\beta \right] = \int_{-\infty}^{+\infty} f(\ln \tau) \exp \left(- \frac{t}{\tau} \right) d(\ln \tau)$$

or related to dynamical variables

$$\frac{I(\mathbf{Q}, t)}{I(\mathbf{Q}, 0)} = \exp [- (\Gamma t)^{2/3}]$$

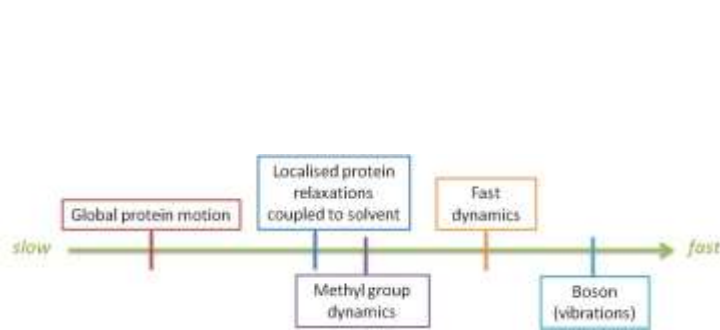
- Relaxation times
- Measure of heterogeneity
- Fraction moving
- Geometry
- Span of motion

1. Protein Dynamics

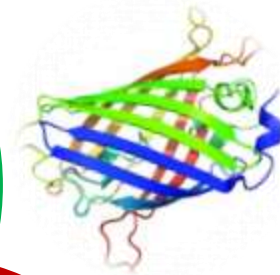
“The protein molecule model resulting from the X-ray crystallographic observations is a platonic protein, well removed in its perfection from the kicking and screaming stochastic molecule that we infer must exist in solution.”



Gregorio Weber
(1916-1997)

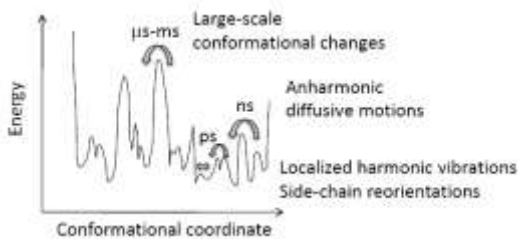
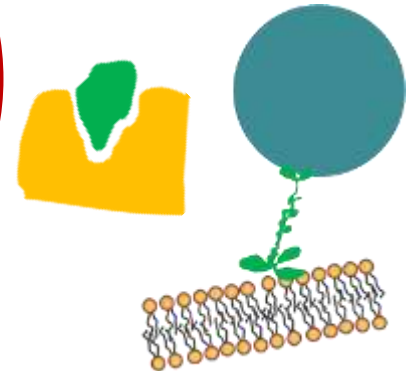


Structure

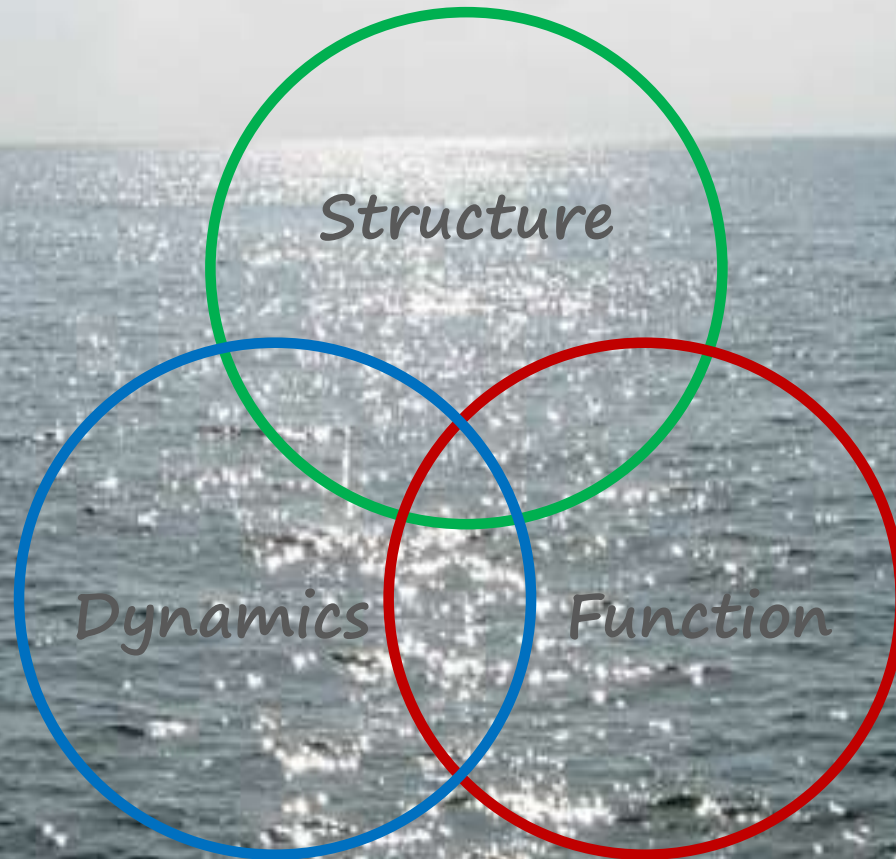


Dynamics

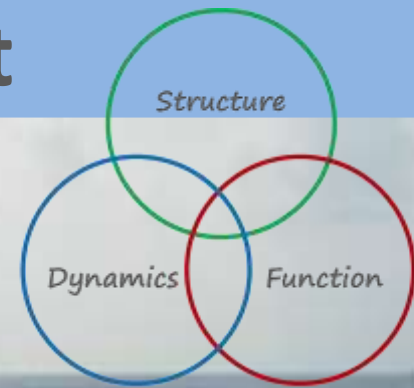
Function



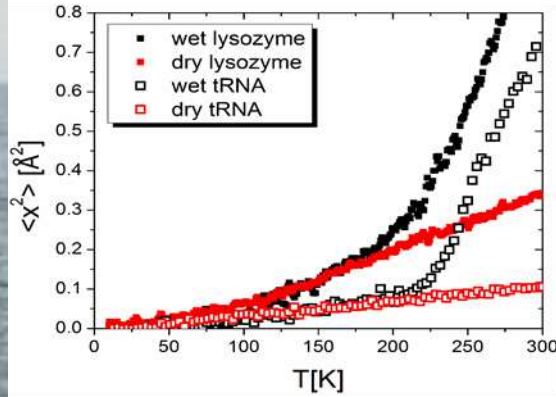
plus the Role of the Solvent



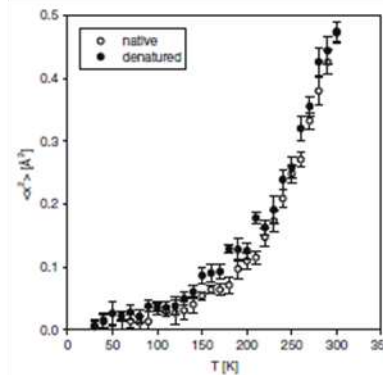
plus the Role of the Solvent



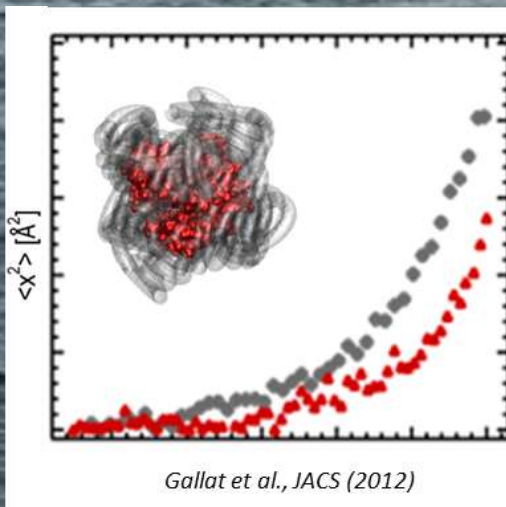
The dynamical transition



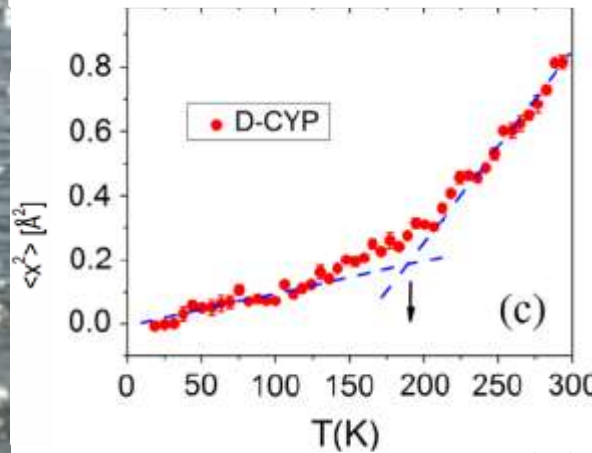
Caliskan et al, JACS (2006)



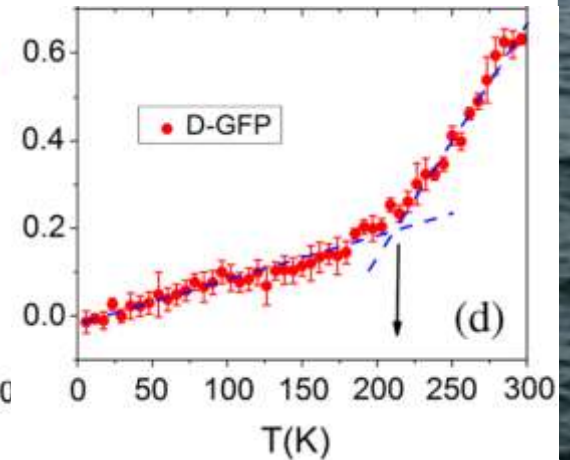
Mamontov et al, JBioPhys (2010)



Gallat et al., JACS (2012)



(c)



(d)

Hong et al., Phys Rev Lett (2017)

Solvent and Bio/Cryopreservation



Trehalose

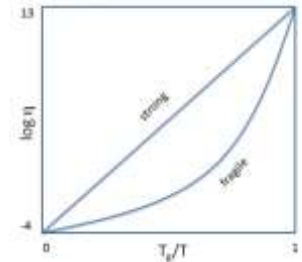


$T_g \approx 390\text{K}$
 $10^{15} \text{ mPa}\cdot\text{s}$
Fragile

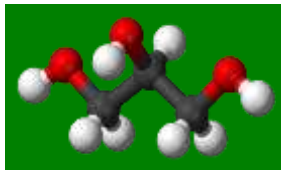
- What is the **best solvent**?
- Understanding **stabilization mechanism**
- Roles of viscosity, packing and fragility



Lysozyme



Glycerol



$T_g \approx 190\text{K}$
 $1200 \text{ mPa}\cdot\text{s}$
Strong

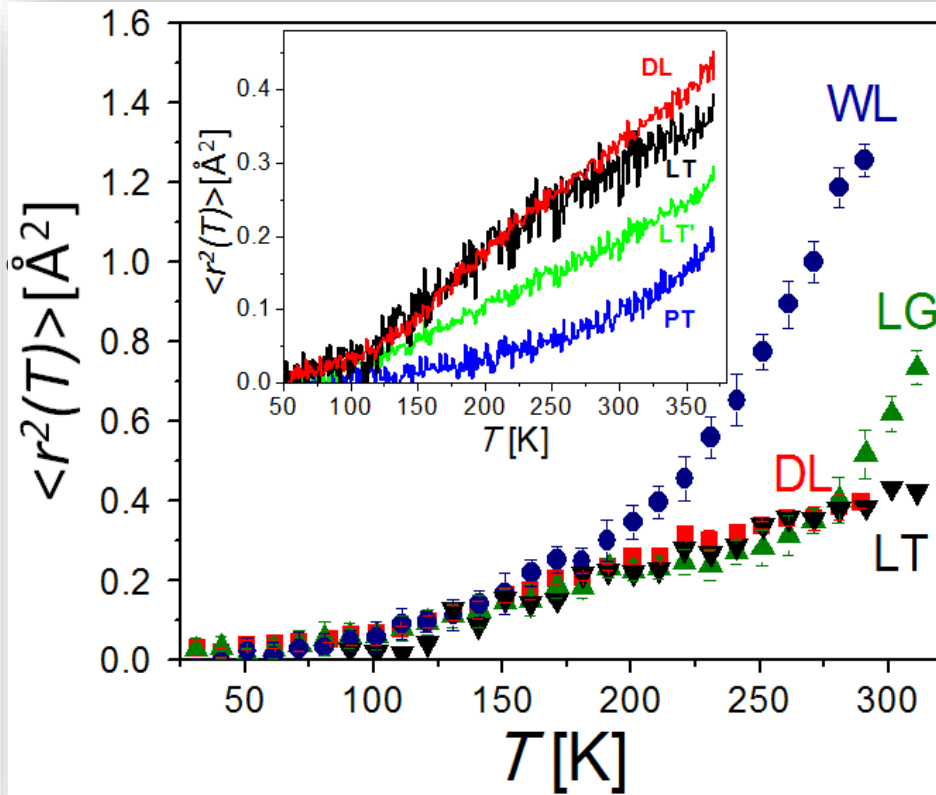
Water
 $1 \text{ mPa}\cdot\text{s}$



- Hydrated powders 1:1
- d-solvents
- HFBS backscattering (0.1-4ns) @ NIST
- Disk Chopper Spectrometer (0.01-50ps)



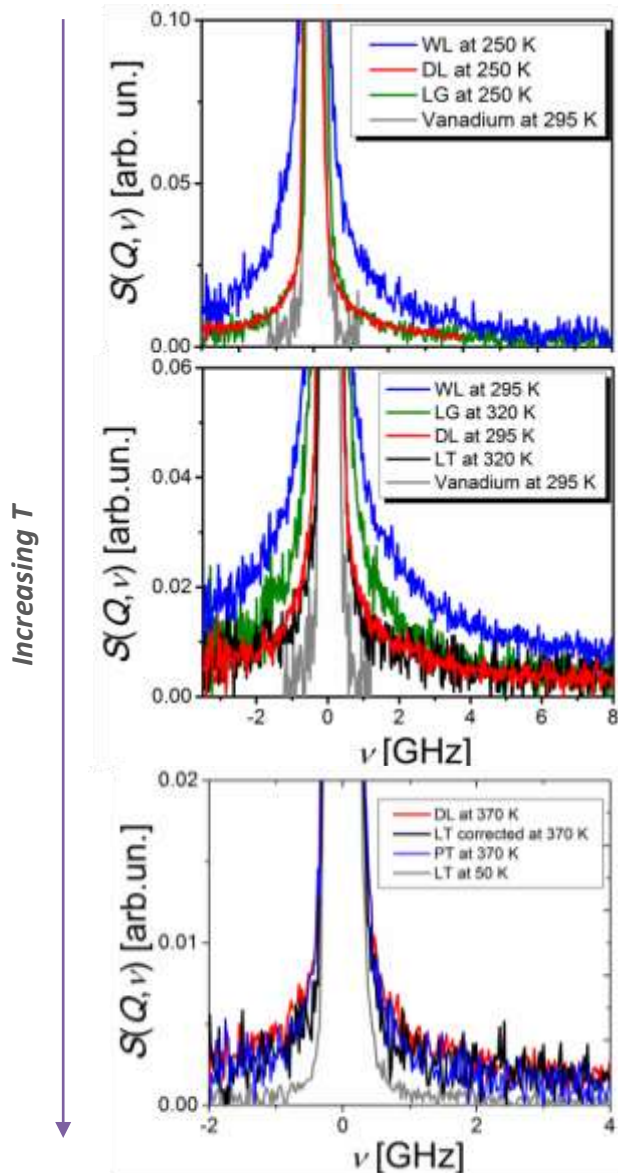
Mean-squared displacements



- CH_3 transition at 100K
- dynamical transition at $\sim 200\text{K}$ for hydrated Lyz (WL)
- dynamical transition at $\sim 270\text{K}$ for Lyz in glycerol (LG)
- no transition for Lyz in trehalose (LT)
- lower mobility in LT vs. dry Lyz (DL)

$$\langle r^2(T) \rangle = -3Q^{-2} \ln[I_{el}(Q, T)/I_{el}(Q, 10K)]$$

Nanosecond dynamics

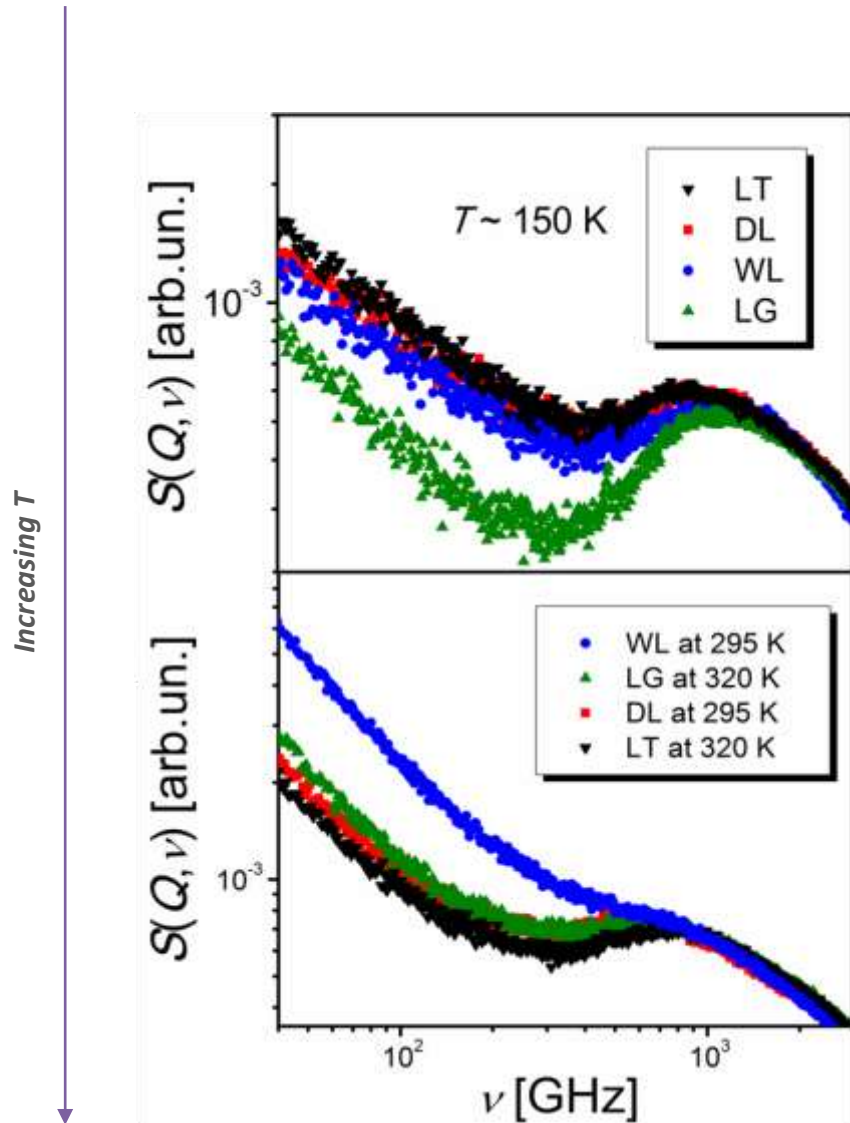


- Localised motions in DL and LG, unaffected by solvent
- water allows further flexibility

- at RT increased mobility in DL, LT
- LG increased flexibility, 50K above Td
- water strongest, 100K above Td

- Localised motions in DL out of energy window
- QENS broadening in LT < in DL

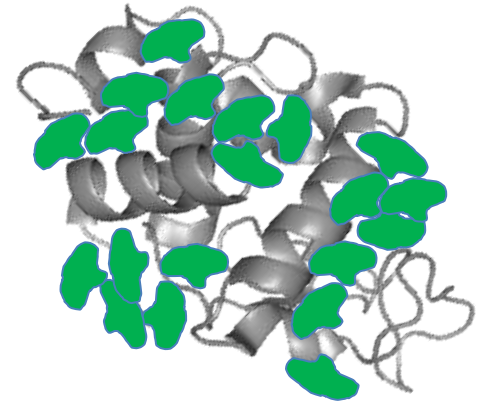
Picosecond dynamics



- Fast conformational transitions and Boson peak at 1THz
 - DL and LT have highest mobility
 - LG has strongly suppressed dynamics
-
- WL shows strongest dynamics
 - LG slightly enhanced than LT
 - LT weaker than DL even at lower T

Emerging picture for **low temperatures**

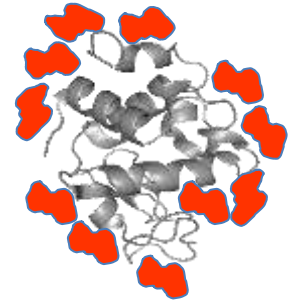
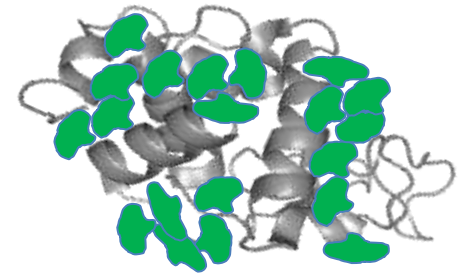
- Small **glycerol** molecules form rigid-well packed structures with weak fluctuations; can interact very efficiently with protein surface
- Reduced cage sizes – also for water
- Strong glass
- Higher viscosity compared to water



- Larger cages and less efficiently packed; poor interaction with protein
- **Trehalose** fast conformational fluctuations
- Trehalose softens dynamics

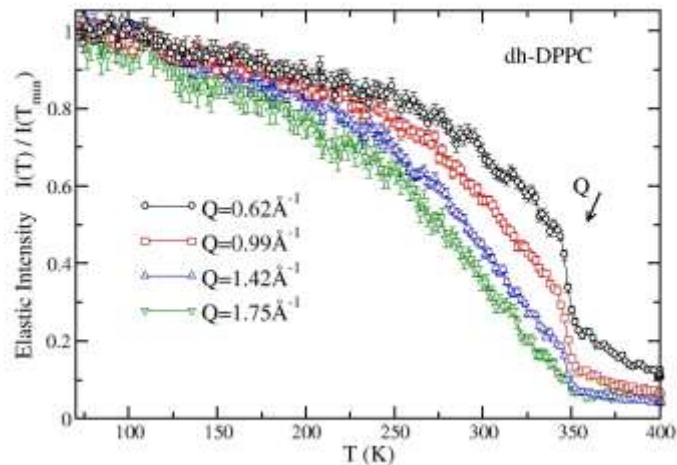
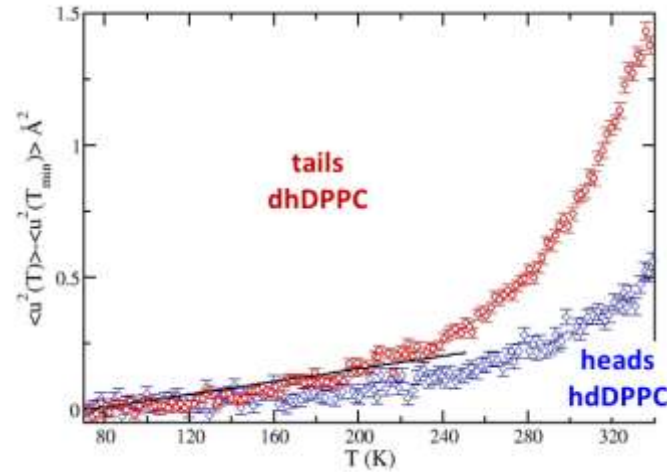
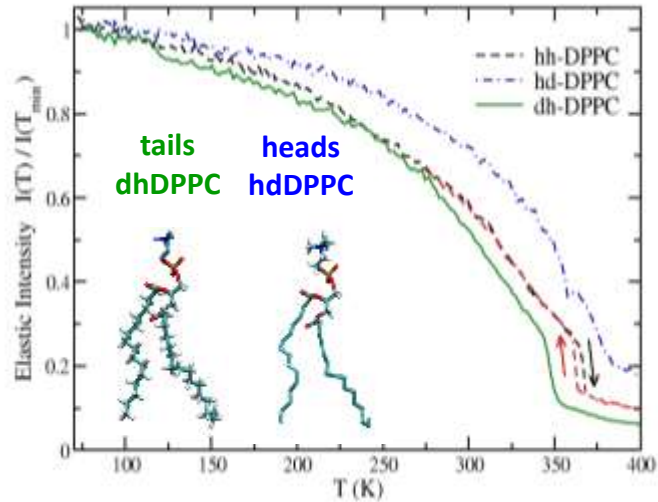
Emerging picture for high temperatures

- Viscosity and high T_g of trehalose controls dynamics
- Water and glycerol are above T_g thus promote dynamics
- Above RT, trehalose suppresses onset of denaturation and unfolding, due to its high T_g
- Suggests that trehalose forms a glassy 'crust' surrounding protein
- Similar observation as for lipids vesicles



- For low temperatures, ie. for cryopreservation, glycerol is the best choice
- At RT and above, for biopreservation and stability, trehalose is the best choice
- Complex interplay between packing and fragility

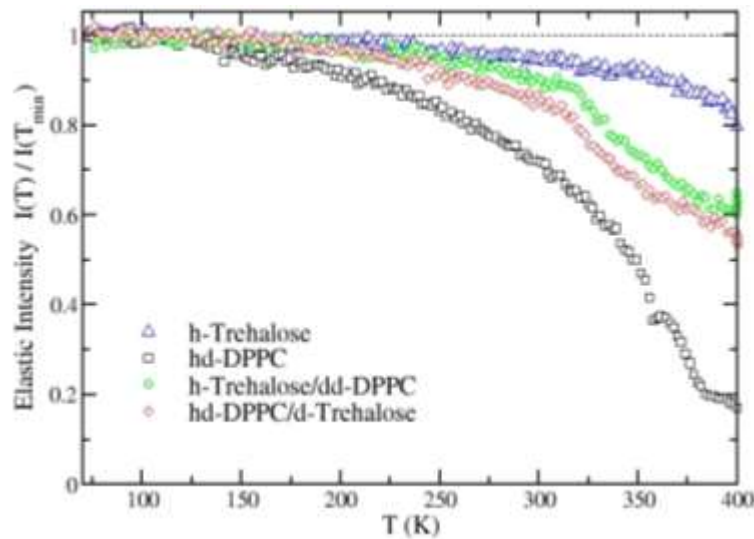
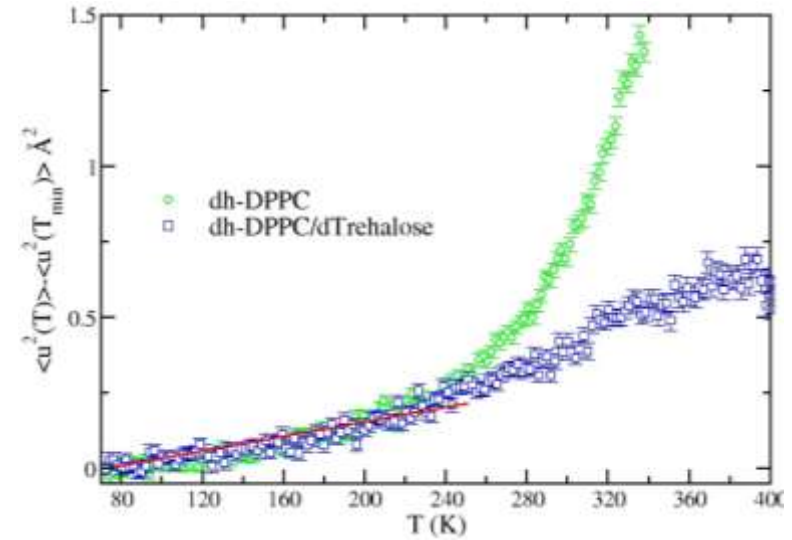
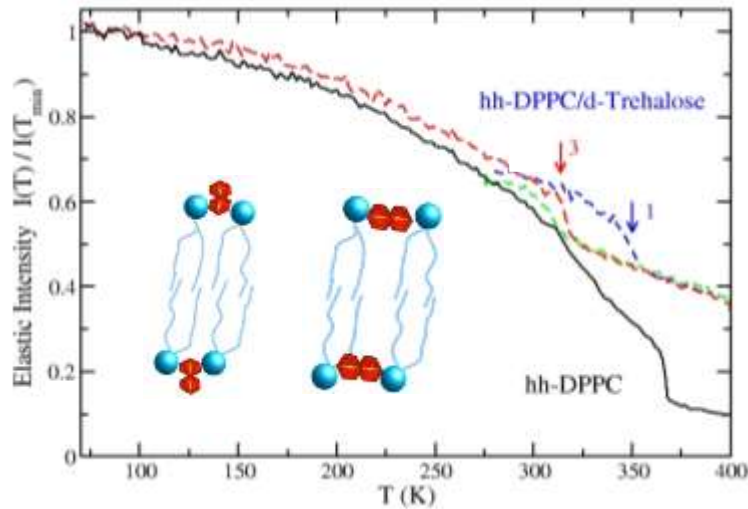
2. Molecular View of Lipid Melting



Sample	DSC	QENS
Protonated lipid (hh-DPPC)	378 K	363–367 K
Heads deuterated lipid (dh-DPPC)	366 K	—
Tails deuterated lipid (hd-DPPC)	367 K	348 K

- Tails responsible for T_m
- Heads less mobile than tails
- Spatial dependence of tail melting

Incorporation of trehalose



- Trehalose decreases T_m
- Partial incorporation, then full
- Strong association to heads
- Significant decrease in mobility

Quantifying Mobility

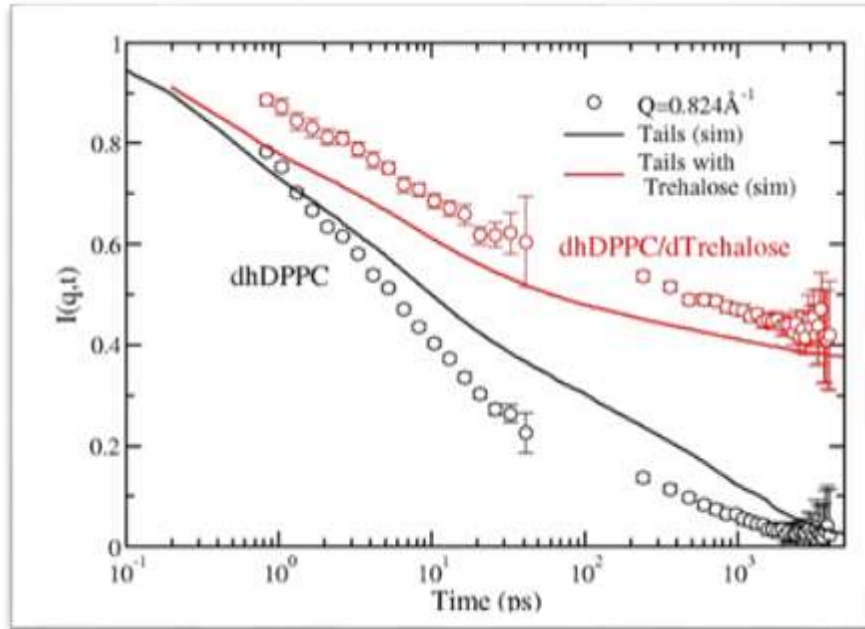


TABLE 3 Parameters of best fit for the model of diffusion inside a sphere using a sum of spheres distributed linearly along the lipidic tails ($n = 2, 3, \dots, 16$)

	Temperature (K)	R_{min} (Å)	R_{max} (Å)
Experiments	310	0.14 ± 0.05	3.20 ± 0.09
	330	0.47 ± 0.12	5.59 ± 0.24
	350	0.79 ± 0.03	5.85 ± 0.16
	370	0.86 ± 0.08	6.19 ± 0.45
	395	1.15 ± 0.07	6.44 ± 0.37
	308*	0.18 ± 0.08	2.75 ± 0.10
	333*	0.72 ± 0.09	4.44 ± 0.20
	370 [†]	1.3	4.2
	390 [†]	1.6	4.7
	Simulations	310	0.45 ± 0.02
330		0.54 ± 0.06	3.91 ± 0.17
350		0.71 ± 0.05	3.83 ± 0.08
370		0.66 ± 0.07	4.56 ± 0.23
395		0.77 ± 0.11	4.86 ± 0.23

Results from literature QENS studies probing alkyl tail mobility are also shown.

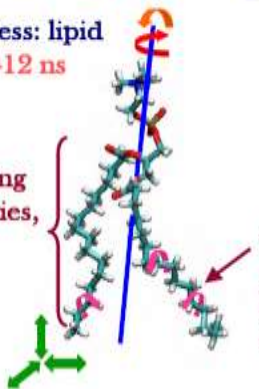
*Oriented multilayers of fully protonated DPPC with 12 wt % water. Data from König et al. (10).

[†]Dicopper alkanolate complexes with deuterated terminal methyl groups (R_{max} corresponds to $n = 15$). Data from Carpentier et al. (43).

Fast process: trans/gauche isomerization 7-45 ps

Slow process: lipid rotation 2-12 ns

If varying mobilities, $\beta < 1$



$$I(q,t) = A K W W_i K W W_f$$

$$K W W_f = E_i + (1 - E_i) \exp\left[-t / \tau_i\right]^\beta$$

- > τ : characteristic time
- > β : distribution of times
- > E : motion in restricted geometry

τ : typical time scale for trans/gauche isomerization

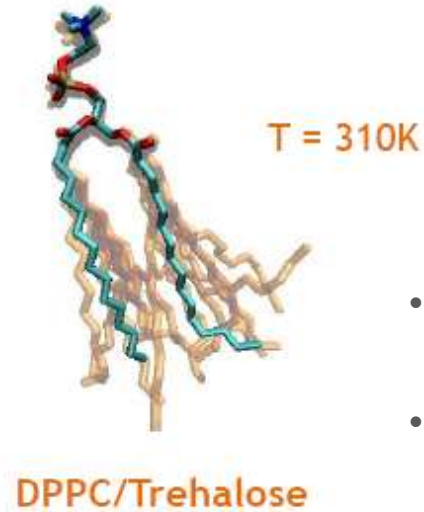
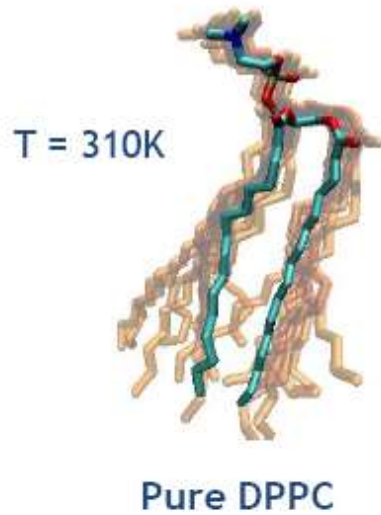
Diffusion in a sphere where r varies linearly with tail position

$$A(Q) = \frac{1}{N} \sum_{r=1}^N \left[\frac{3j_1(QR_r)}{QR_r} \right]^2$$

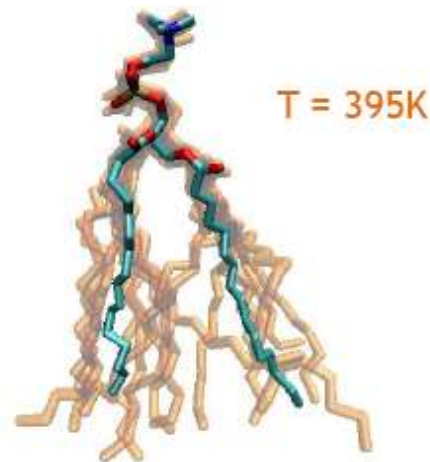
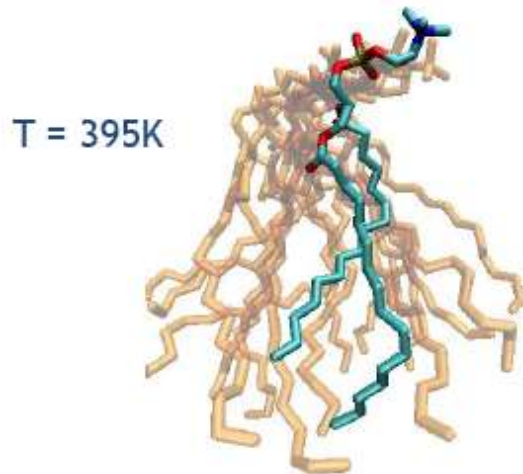
$$R_n = \frac{n-2}{N-2} [R_N - R_2] + R_2$$

Molecular simulations are in good agreement over all lengths and times

Picture of Interaction with Trehalose



- Trehalose decelerates fast conformational transitions along all tail bonds
- Rotations or any slower motions are significantly suppressed
- Spatial extent is significantly decreased



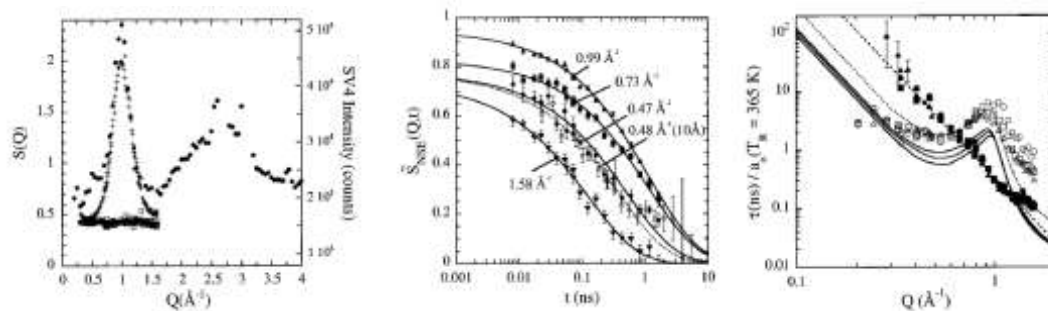
Collective dynamics - coherent

Collective Dynamics – Coherent QENS

Let's consider the **coherent part of the correlation function**. Interpretation is trickier and you need to take into account the **structure factor, $S(Q)$** , ie. how atoms are distributed in space. Typically longer times and measurements in real time (not energy).

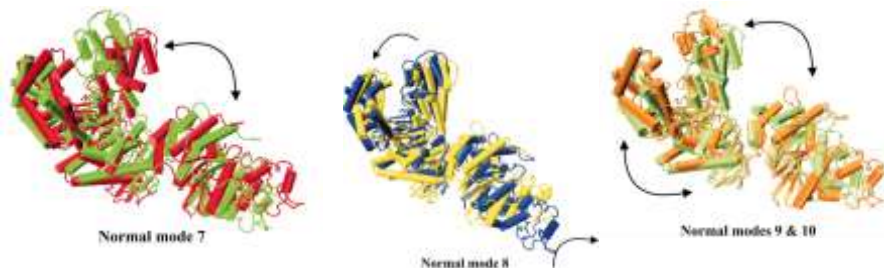
$$I_{\text{coll}}(Q, t) \approx I_{\text{self}}(Q, t) \left(\frac{Q}{\sqrt{S(Q)}}, t \right)$$

$$\tau_{\text{coll}}(Q, T) \approx a(T) \tau_{\text{self}}(Q, T) S(Q)^{1/\beta}$$



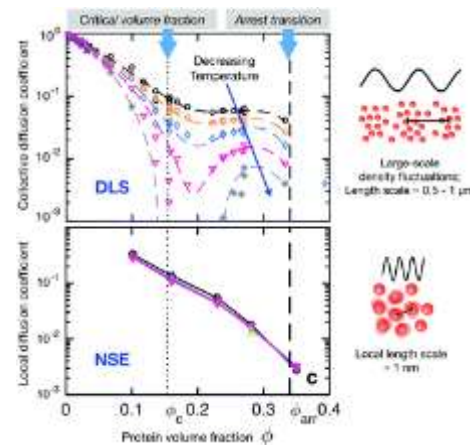
Farago et al, PRE, (2002)

Assignment of normal modes of proteins in solution



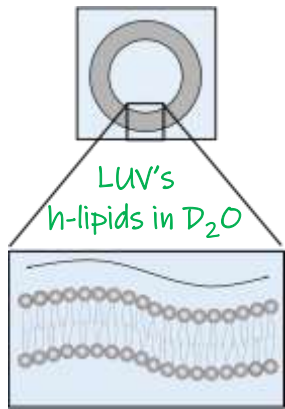
Bu et al, PNAS (2005)

Collective diffusion of proteins in solution



Stradner and Schurtenberger, Soft Matter (2019)

3. Bending Elasticity in Model Membranes



Zilman-Granek theory (*PRL* 77, 4788, 1996)

$$\frac{I(q, t)}{I(q, 0)} = \exp[-(\Gamma t)^{2/3}] \quad \dots \text{measured by NSE}$$

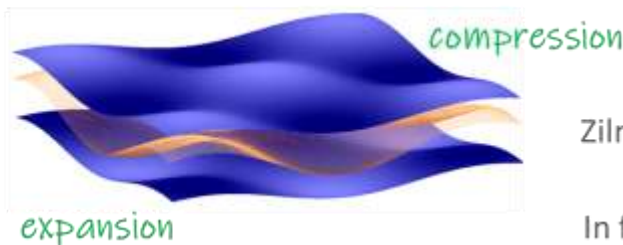
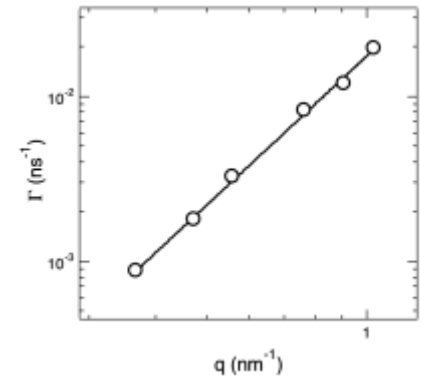
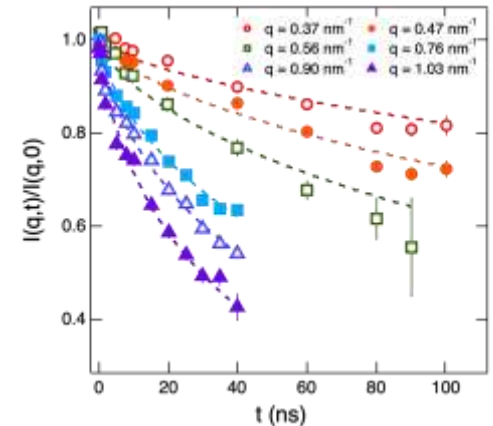
$$\Gamma_{bend} = 0.025 \sqrt{\frac{k_B T}{\kappa} \frac{k_B T}{\eta}} q^3$$

Bending modulus
Solvent viscosity

$$\tilde{\kappa} = \kappa + d^2 \kappa_A$$

Effective bending modulus
Compressibility modulus

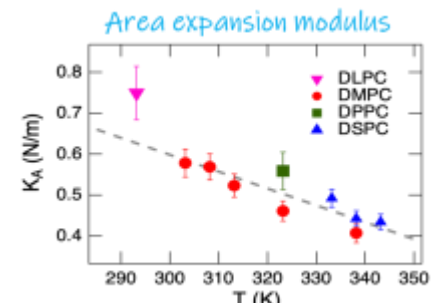
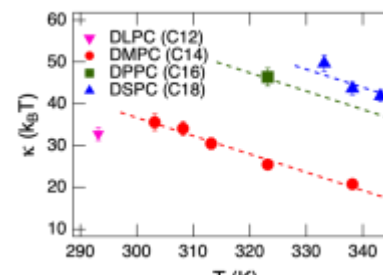
Watson and Brown,
Biophys. J., 98
(2010).



Zilman-Granek theory is for a homogeneous thin sheet

In fact, dissipation of stored energy through solvent and membrane

Additional motion controlled by K_A and both solvent and membrane viscosities

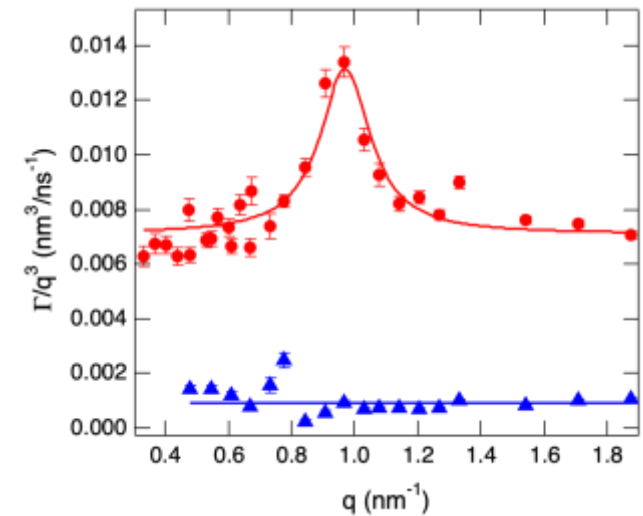
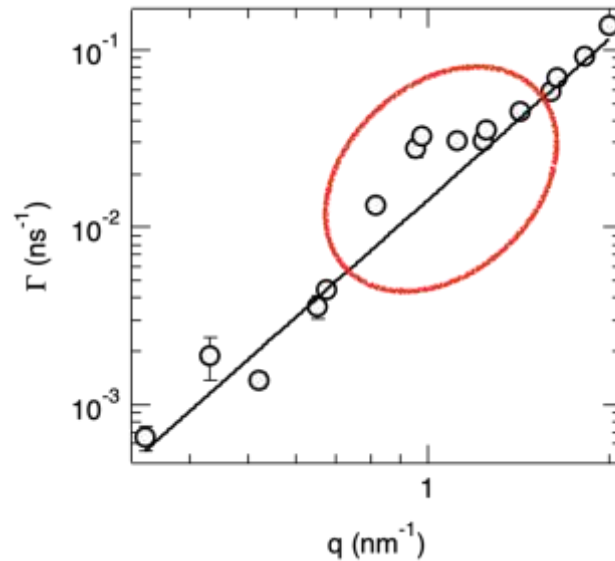
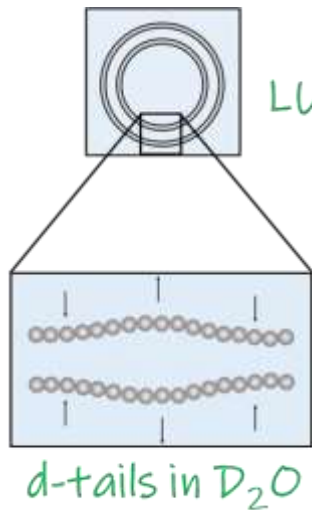


$$K_A = \frac{\beta \kappa}{h_c^2}$$

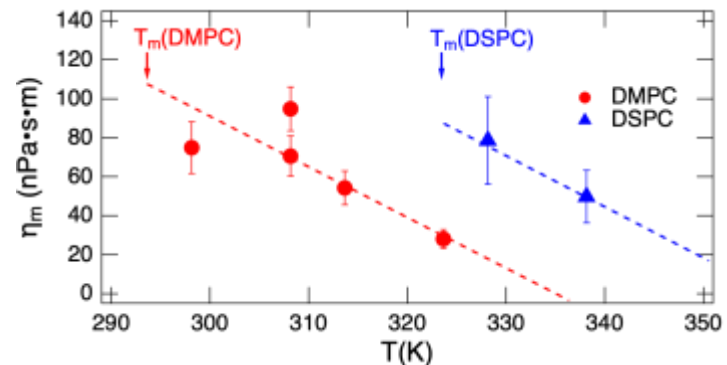
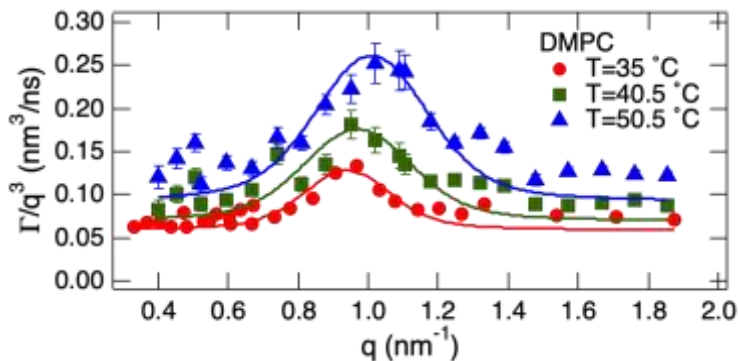
Monolayer constant
 $\beta = 12$ fully couple (like a slab)
 $\beta = 48$ fully uncoupled (each leaflet can slide past each other)
 $\beta = 24$ polymer brush model for lipids

Courtesy of M. Nagao

Bilayer thickness fluctuations



Excess dynamics observed at bilayer thickness only visible in fluid phase + estimation of the membrane viscosity



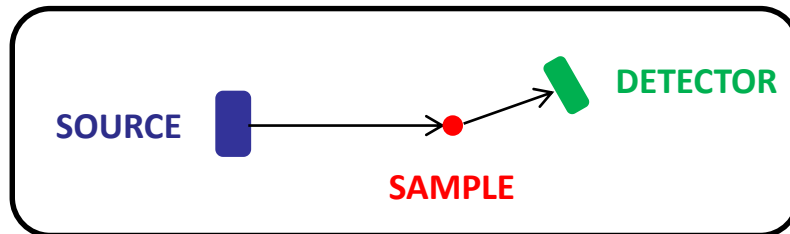
$$\frac{\Gamma}{q^3} = 0.0069 \sqrt{\frac{k_B T}{\kappa}} \frac{k_B T}{\eta} + \frac{K_A k_B T}{\eta_m q_0^3 k_B T + 4 \eta_m q_0 K_A A_L (q - q_0)^2}$$

Instruments

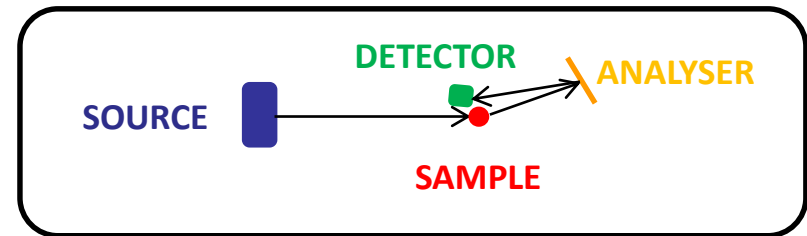
Neutron Spectrometers

Frequency domain $S(Q, \omega)$

Direct geometry



Indirect geometry or Backscattering



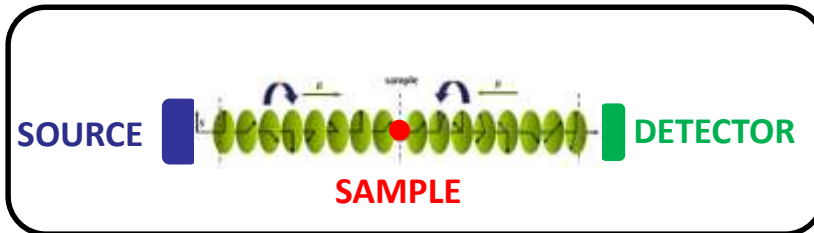
- Lower energy resolution
- Larger ΔE and Q range
- Flexible in choosing Q - E space
- Repetition Rate Multiplication
- picosecond motions (diffusion, rotations, trans-gauche conformations)
- Self-correlation function (H)
- IN5, LET, CNCS ...

- Higher energy resolution
- Smaller ΔE and Q range
- Fixed Q - E space
- Picosecond-nanosecond (diffusion, rotations, trans-gauche conformations)
- Self-correlation function (H)
- DNA, IRIS, IN16B, HFBS

Neutron Spectrometers

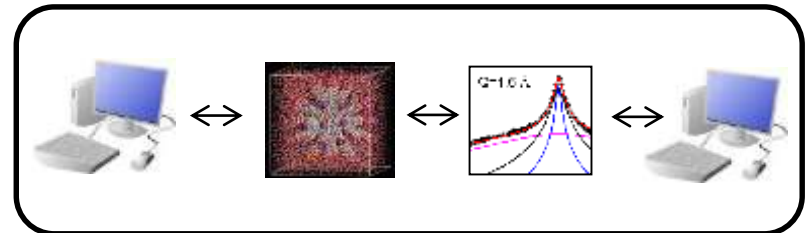
Time domain $I(Q,t)$

Neutron Spin Echo



- Highest energy resolution
- Typically lower Q range
- Nanosecond (normal modes, bending elasticity, fluctuations)
- Collective dynamics (H/D)
- IN15, NIST-NSE

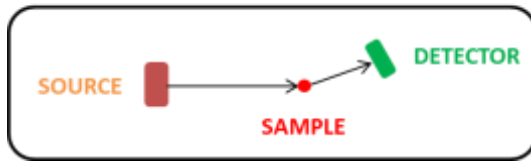
Molecular Dynamics Simulations



- Flexible parameter space
- Large parameter space
- Need to optimize and choose FF
- Directly comparable with neutron data

Time-of-flight Spectrometers

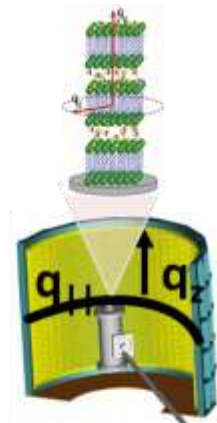
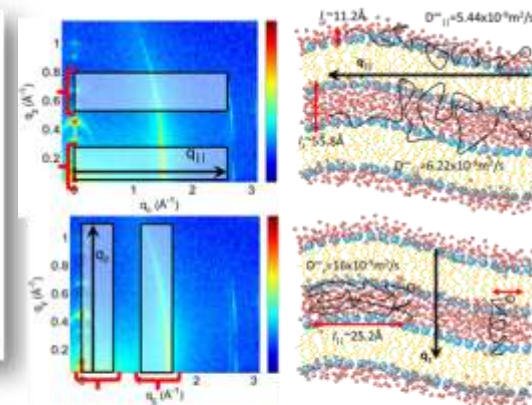
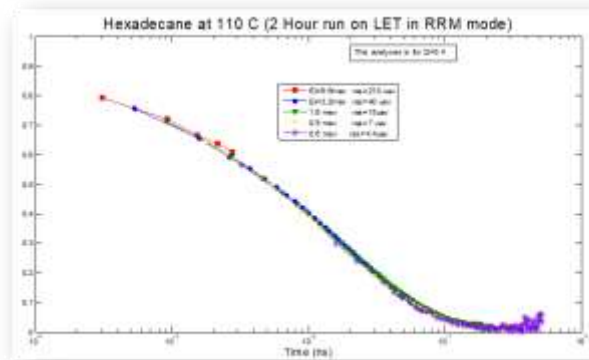
Direct geometry



- Resolution 10-100 μ eV, but flexible Q and E range.
- Measure picosecond motions and fast collective modes
- Good choice for QENS over a broad range, especially making use of rep-rate multiplication
- Position sensitive detectors enable 2D studies
- Recently advances with polarisation analysis enables separation of coherent/incoherent scattering

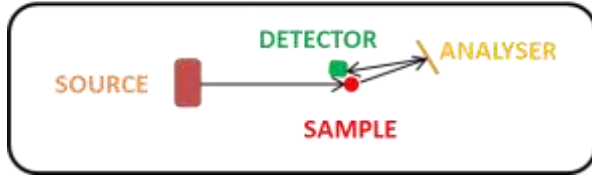


LET on ISIS



Backscattering Spectrometers

Indirect geometry or Backscattering



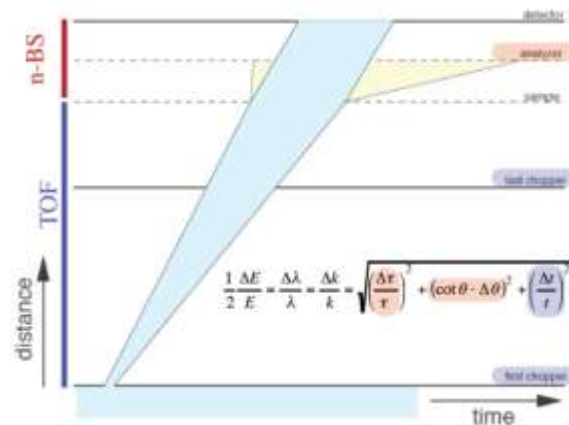
$$n\lambda = 2d\sin\theta$$

$$\frac{1}{2} \frac{\Delta E}{E} = \frac{\Delta\lambda}{\lambda} \sim \frac{\delta d}{d} + \cot(\theta)\delta\theta$$



HFBS @ NIST

- Energy resolution (ueV) selected by analyser crystals. Changing crystal changes δE , ΔE and ΔQ .
- Measure pico-nanosecond motions.
- Reactor ~ 0.9 u eV, limited energy range, fixed Q-range.
- Spallation ~ 3 -20ueV, broader/flexible ΔE but similar Q range
- Good choice for elastic fixed window scans and QENS



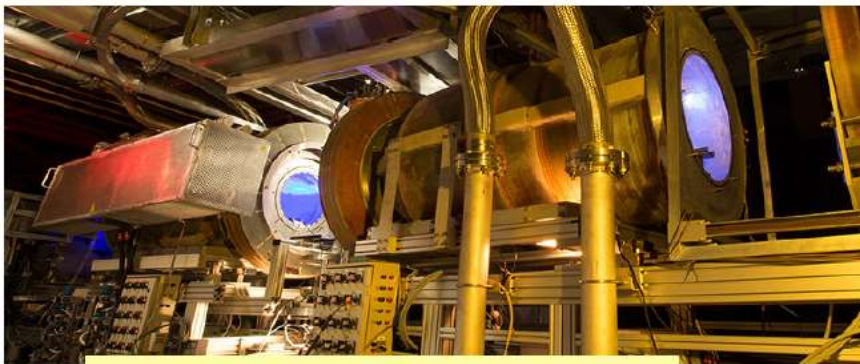
BASIS @ SNS



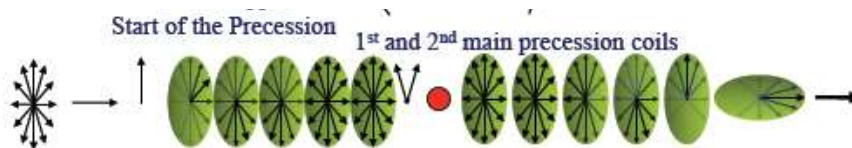
Neutron Spin-Echo

- NSE uses the neutron's **spin polarisation** to encode the difference in energies between incident and scattered beams. Very high energy resolution.
- Neutrons perform **Larmor precessions** in two antiparallel magnetic fields, before and after the sample, resulting in polarization of the neutrons.
- Precession angles are equal and opposite and the difference is analysed at the detector. Small energy transfers lead to a change in the precession angle and thus a decrease in measured polarization.

$$I_{\text{NSE}}(Q, t) = \frac{I_{\text{coh}}(Q, t) - \frac{1}{3}I_{\text{inc}}(Q, t)}{I_{\text{coh}}(Q, 0) - \frac{1}{3}I_{\text{inc}}(Q, 0)}$$



Today's IN15: measures up to 1 μs



- NSE is the neutron spectroscopy with highest energy resolution
- Time covered is $1\text{ps} < t < 1\mu\text{s}$ (equivalent to neV resolution)
- Momentum transfer range is $0.002 < Q < 4\text{\AA}^{-1}$
- Best signal is at the structural peaks and is very complementary to e.g. SANS

Outline

- Why** do we need to study nanoscale motions in biomolecules
- What is **quasi-elastic neutron scattering**
- What information do **neutrons** give us
- What **instruments** are available
- Examples throughout

Take home messages

1. **Dynamics** provide insights into relevant biomolecules
2. **Instruments** available to probe different motions
3. **Complementary** to structural characterisation
4. Neutrons provide **unique** information